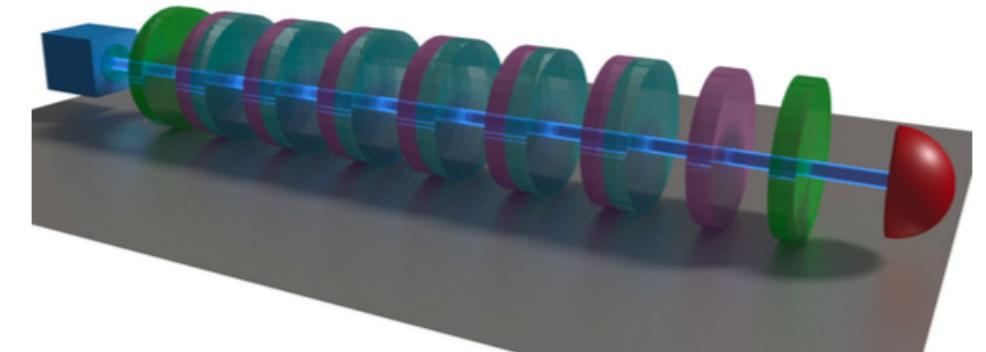
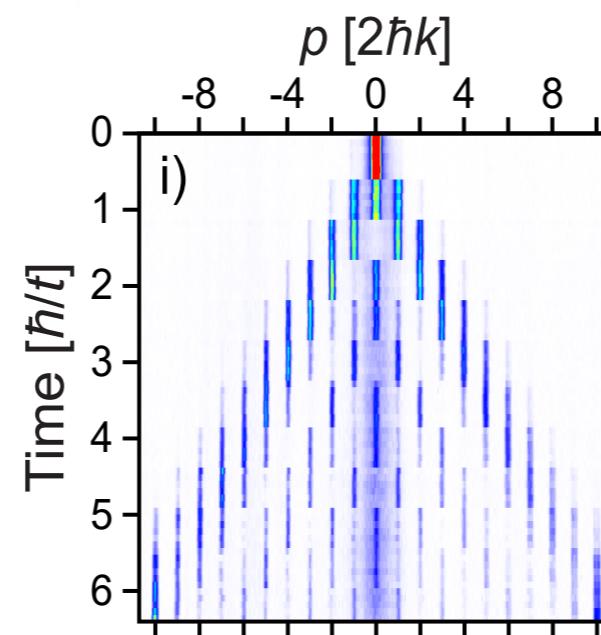
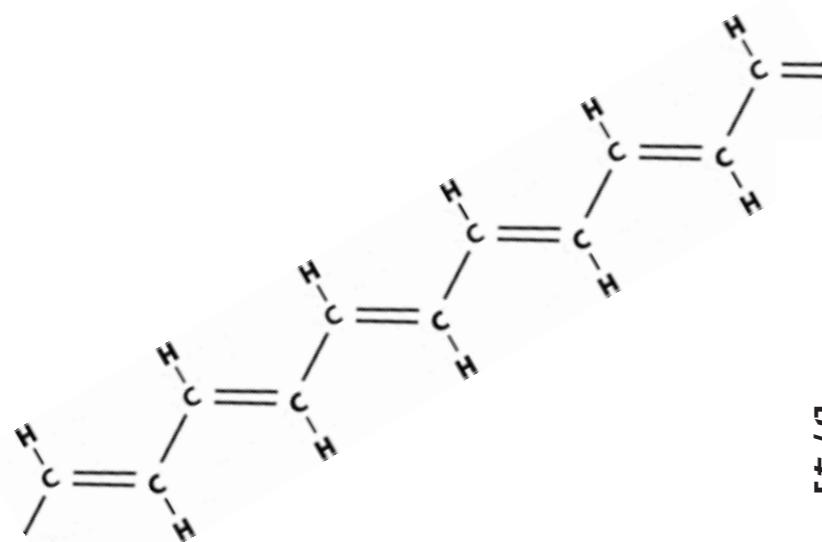


Observation of Topological Anderson & Floquet Insulators

Pietro Massignan



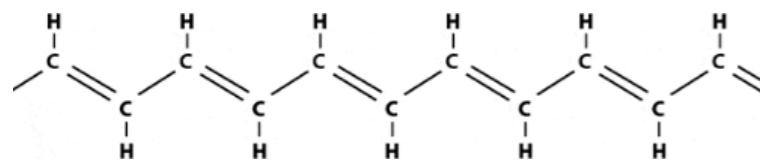
UNIVERSITAT POLITÈCNICA
DE CATALUNYA
BARCELONATECH

 PHYSICS ILLINOIS

SLAM group

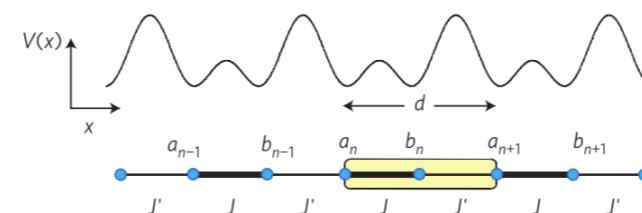
ICFO^R
The Institute of Photonic
Sciences

1D chiral systems



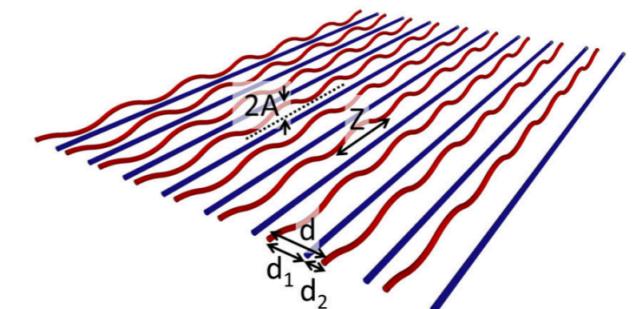
polyacetylene

[Nobel prize in Chemistry 2000]

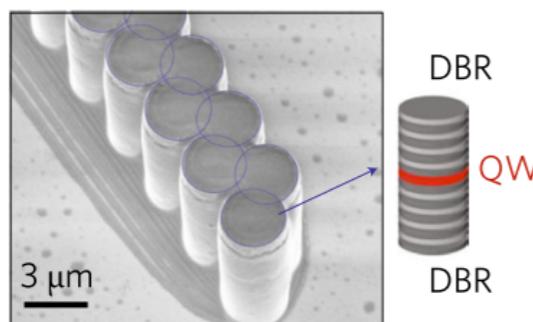


**ultracold atoms
in superlattices**

[M. Atala *et al.*, Nature Phys. 2013]

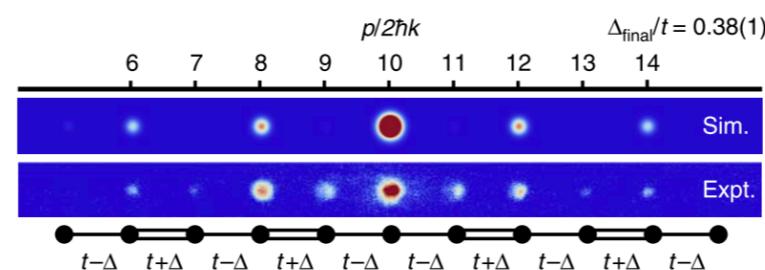


optical waveguides
[Zeuner *et al.*, PRL 2015]



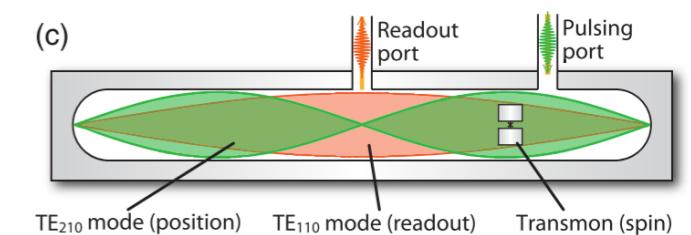
cavity polaritons

[St. Jean *et al.*, Nature Phot. 2017]



**ultracold atoms
in k-space lattices**

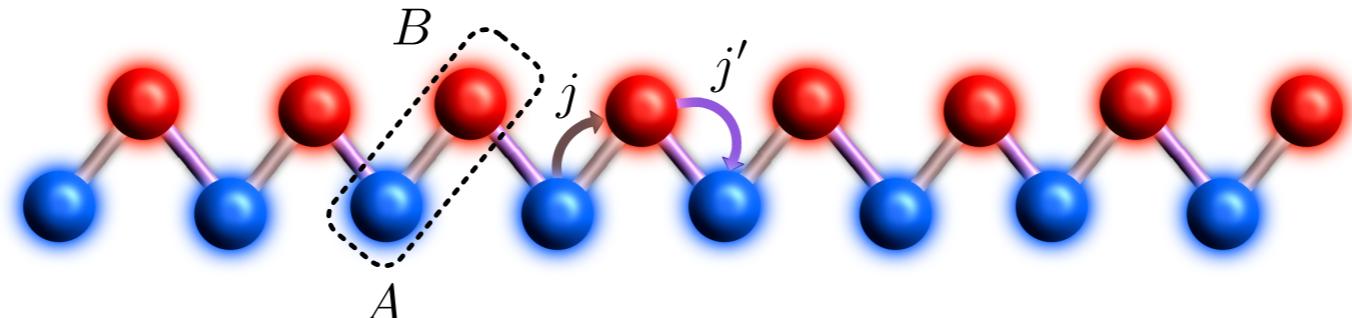
[Meier *et al.*, Nature Comm. 2016]



**SC qubits
in mw-cavities**
[Flurin *et al.*, PRX 2017]

SSH model

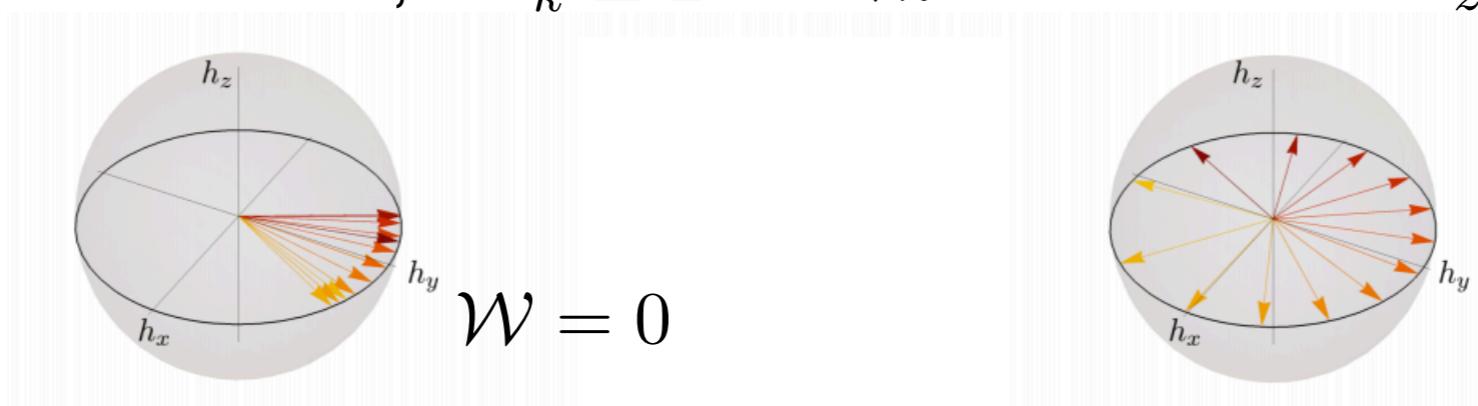
- Spinless fermions with staggered tunnelings:



*Su, Schrieffer & Heeger
Phys. Rev. Lett. (1979)*

*Asbóth, Oroszlány, & Pályi
Lecture Notes in Physics (2016)*

- \exists two sublattices
- \exists a “canonical basis” where H is purely off-diag:
$$H = \begin{pmatrix} 0 & h^\dagger \\ h & 0 \end{pmatrix}$$
- Chiral symmetry: $\Gamma H \Gamma = -H$ (Γ : unitary, Hermitian, local)
- In momentum space: $H_k = E_k \mathbf{n}_k \cdot \boldsymbol{\sigma}$
- In the canonical basis, $\mathbf{n}_k \perp \hat{\mathbf{z}}$ $\forall k$ and $\Gamma = \sigma_z$
- Winding:



The winding W

- \mathcal{W} may be calculated:

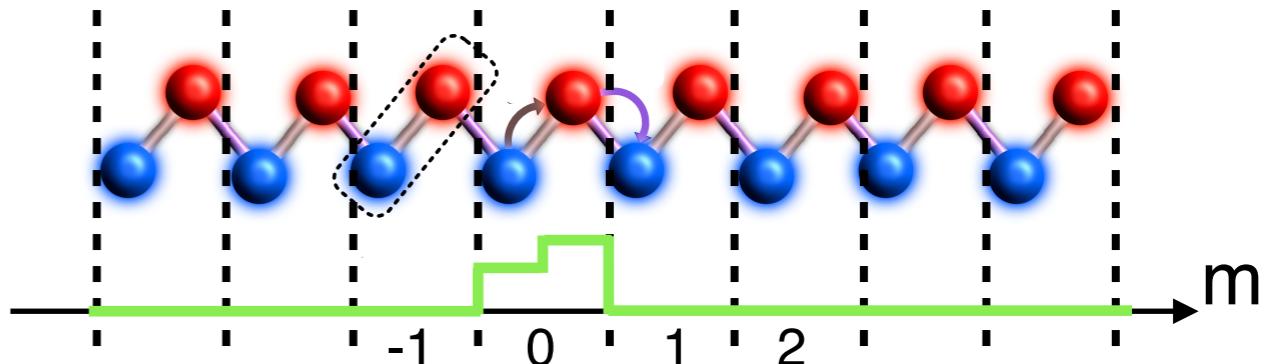
$$H_k = E_k \mathbf{n}_k \cdot \boldsymbol{\sigma}$$

- from \mathbf{n} : $\mathcal{W} = \oint \frac{dk}{2\pi} (\mathbf{n} \times \partial_k \mathbf{n})_z$
 - from the *eigenstates*: $\mathcal{W} = \oint \frac{dk}{\pi} \mathcal{S}, \quad \mathcal{S} = i\langle \psi_+ | \partial_k \psi_- \rangle$
skew polarization
-
- What if the Hamiltonian is not known?
Can one *measure* the winding?

Yes, and it's simple!

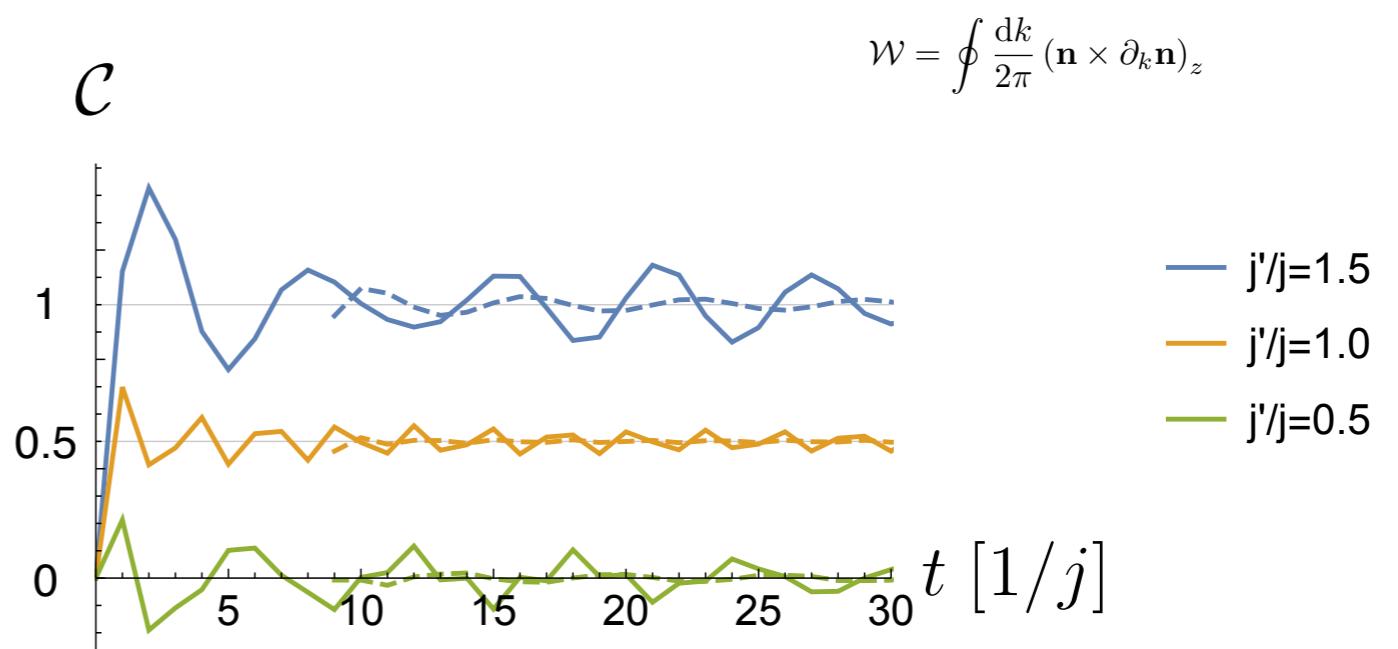
Evolution in real time

- Initial condition
localized on the $m=0$ cell:



- Mean Chiral Displacement:** $\mathcal{C}(t) \equiv 2\langle \widehat{\Gamma m}(t) \rangle = 2[\langle m_A(t) \rangle - \langle m_B(t) \rangle]$

$$\mathcal{C}(t) = 2 \int_{-\pi}^{\pi} \frac{dk}{2\pi} \langle U^{-t} \sigma_z(i\partial_k) U^t \rangle_{\psi_0} = 2 \int_{-\pi}^{\pi} \frac{dk}{2\pi} \sin^2(Et) |\mathbf{n} \times \partial_k \mathbf{n}| \xrightarrow{t \rightarrow \infty} \mathcal{W}$$

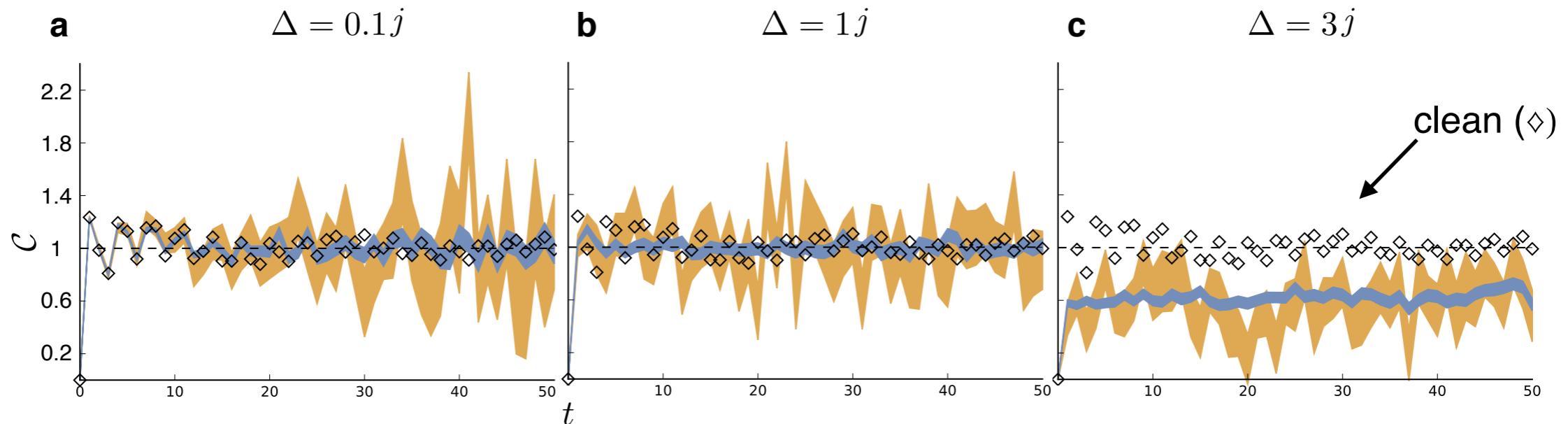


- Bulk measurement*
- Fast convergence

Resistance to disorder

SSH model in the topological phase
+
independent disorder of amplitude Δ on ***all*** tunnelings
+
localized initial condition (randomly-polarized)
+
average over 50 (1000) disorder realizations

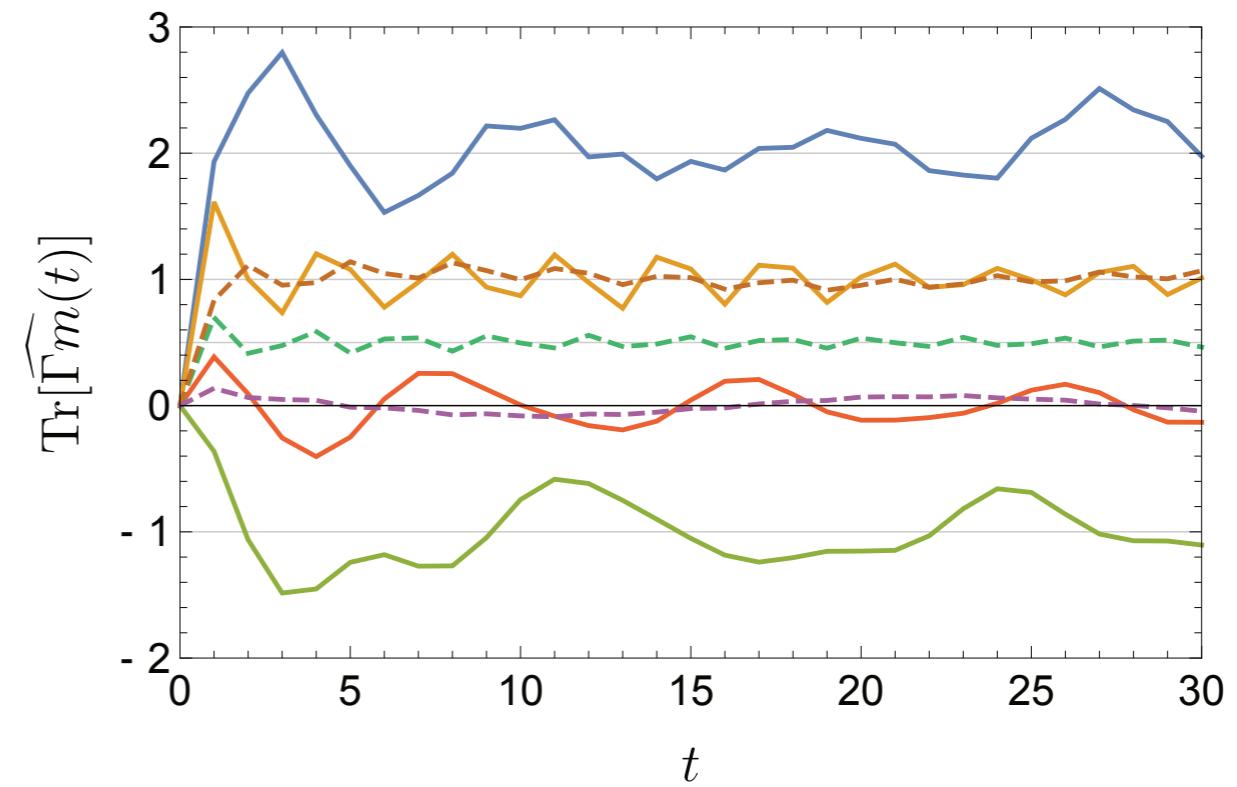
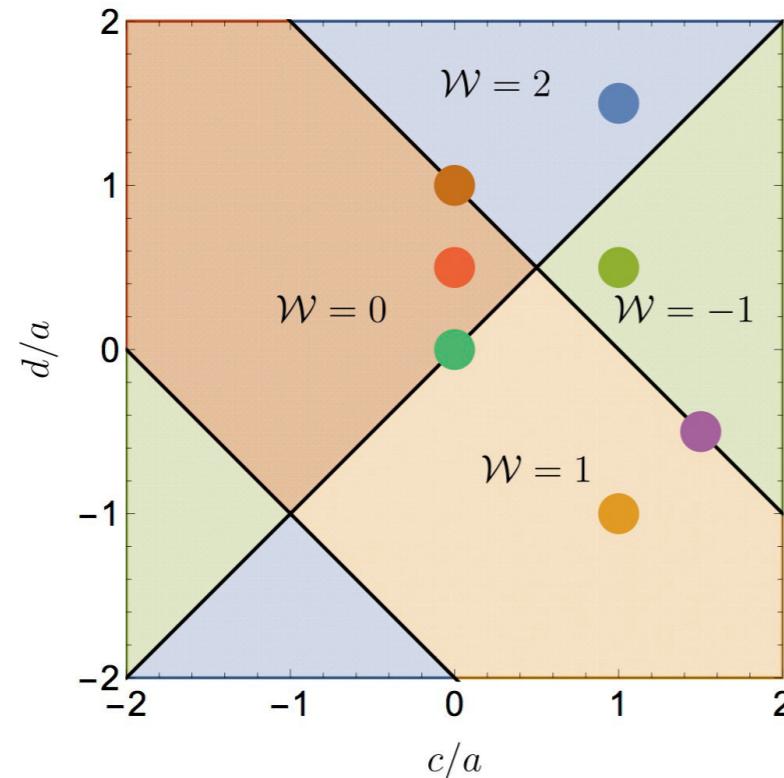
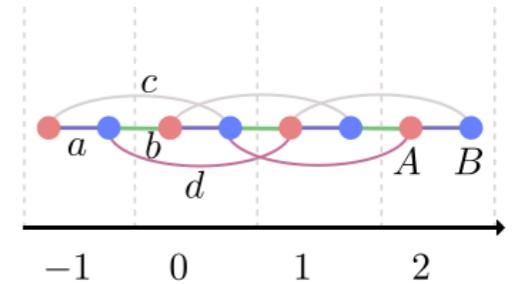
$j' = 2j \rightarrow \begin{cases} \mathcal{W} = 1 \\ \Delta_{\text{gap}} = 2j \end{cases}$



the MCD stays locked to the topological invariant as long as $\Delta < \Delta_{\text{gap}}$

Higher windings

- Extension to long-ranged models:

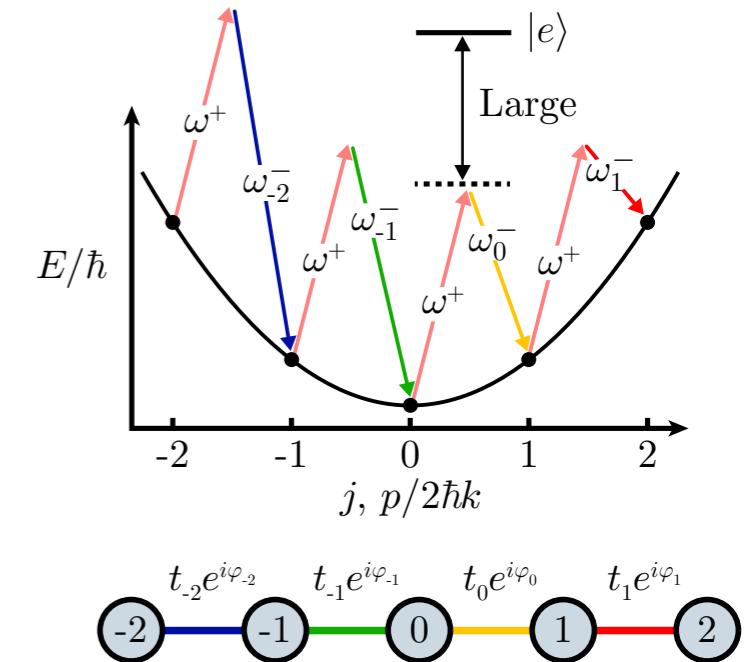
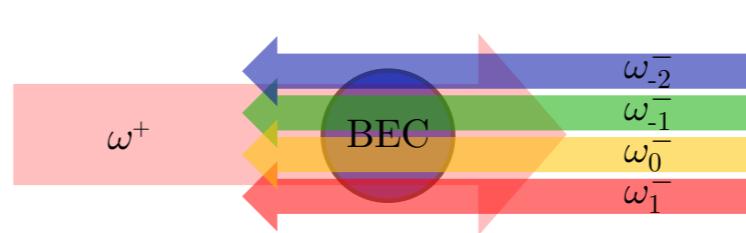


- At critical boundaries: MCD converges to the mean of the winding in the neighboring phases

Atomic wires in Urbana

- Atomic, ~ideal BEC
- Laser-driven coupling of discrete-momentum states

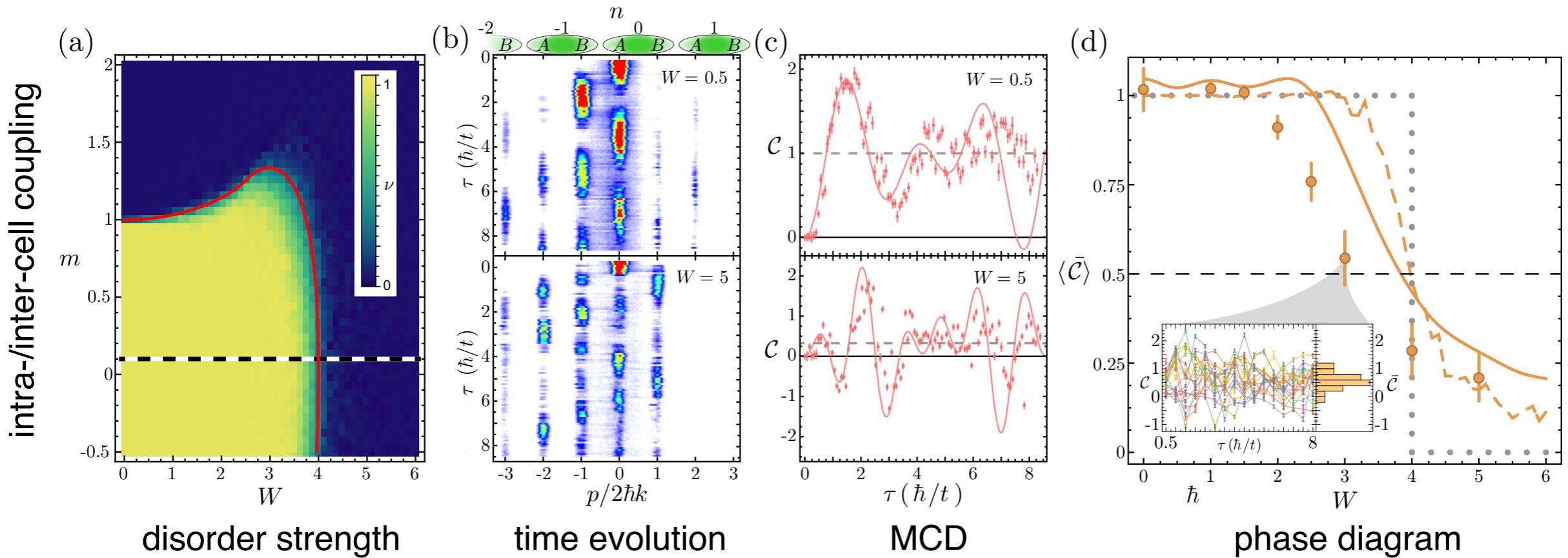
$$H_{\text{eff}} \approx \sum_j t_j (e^{i\varphi_j} |\tilde{\psi}_{j+1}\rangle \langle \tilde{\psi}_j| + \text{h.c.})$$



- 1D Hubbard model with full control on each tunneling strength and phase
- Built-in chiral symmetry

Detecting topology

A topological wire becomes trivial by adding disorder



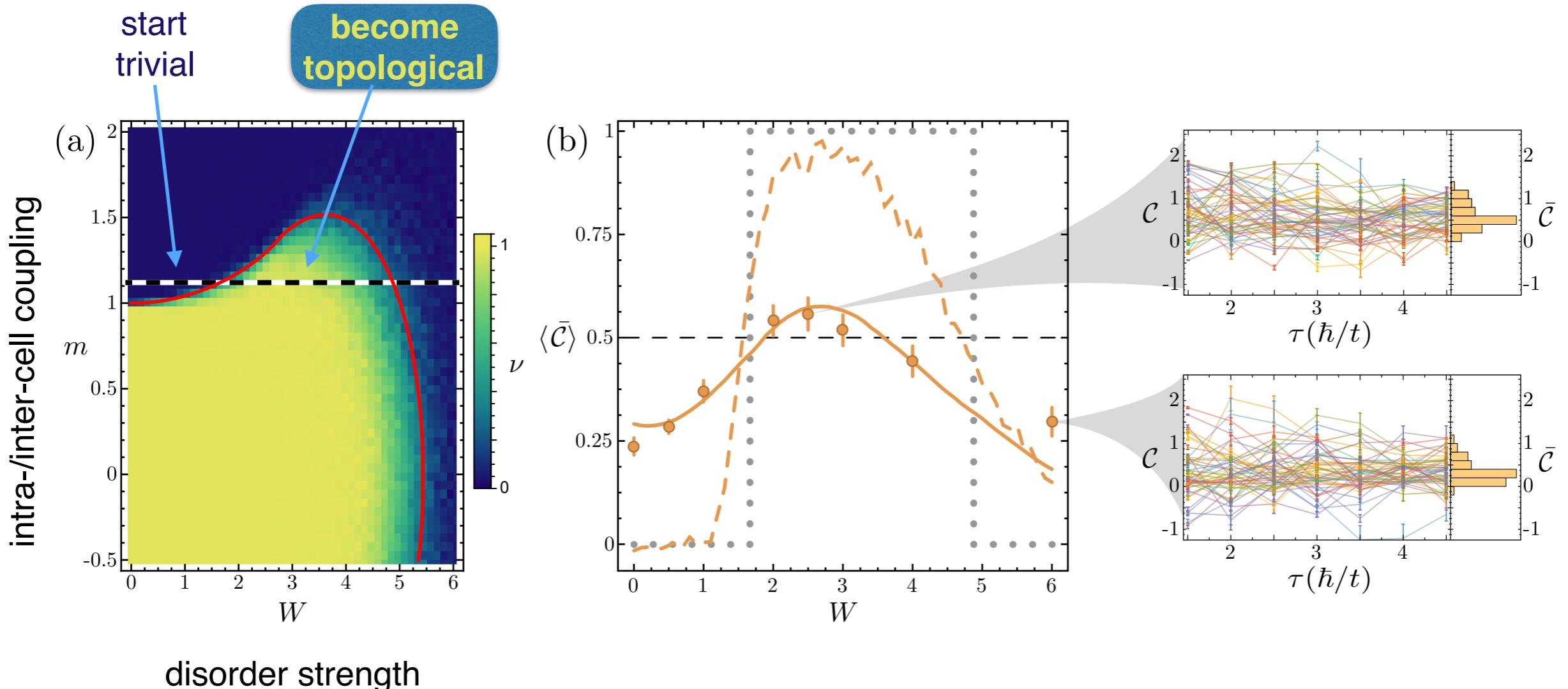
color map: real-space computation of the winding

red line: critical boundary (diverging localization length)

datapoints: experimental measurement of the MCD

Topological Anderson transition

A trivial wire is driven into the topological phase by adding disorder



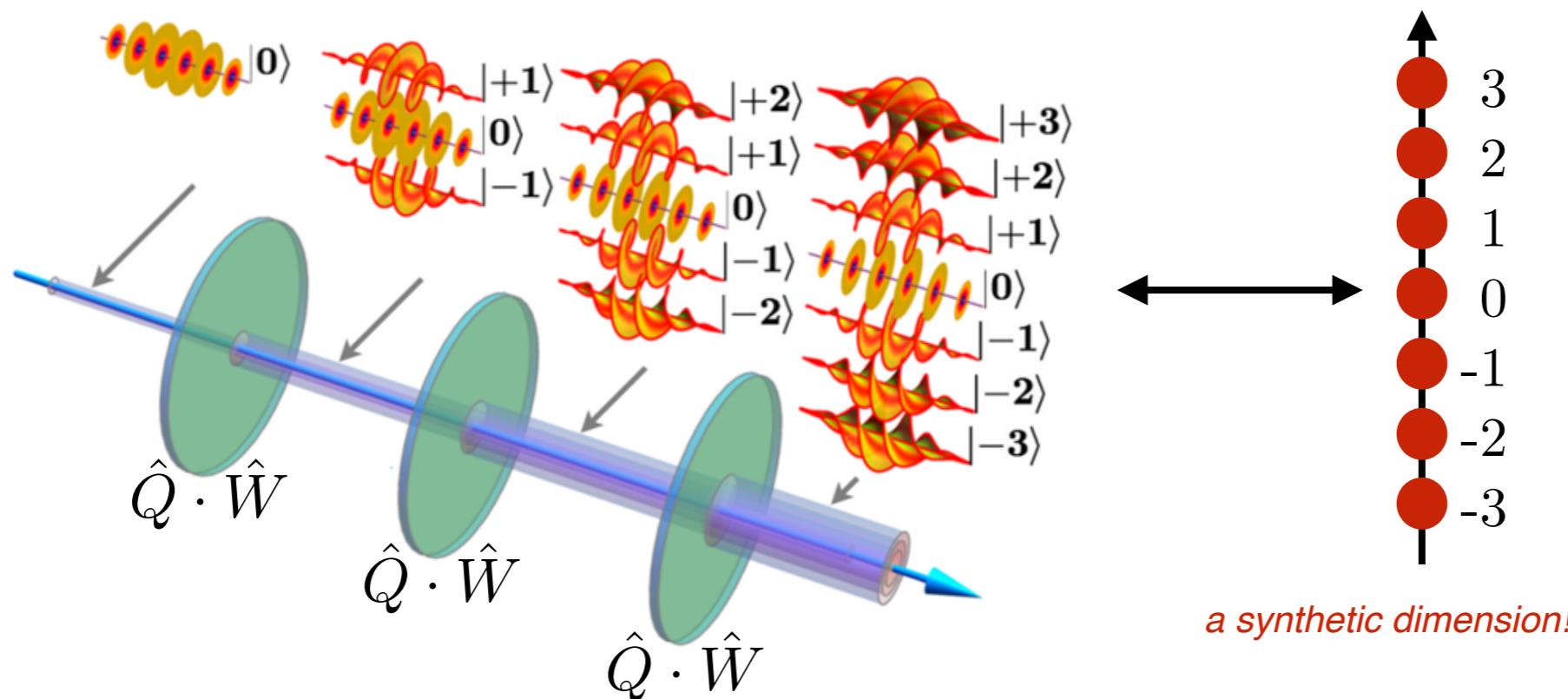
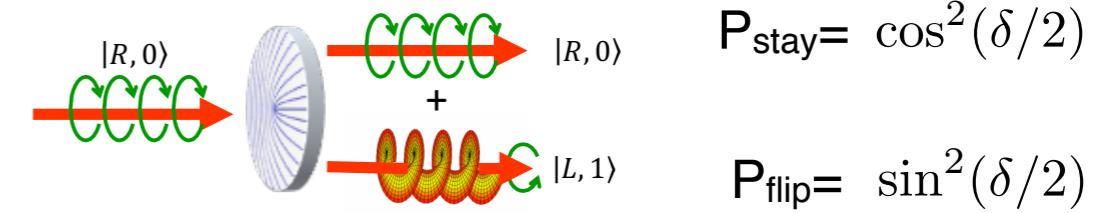
Meier, An, Dauphin, Maffei, PM, Taylor and Gadway,
arXiv:1802.02109

Discrete-Time Quantum Walks with twisted photons

- Cascade of Q-plates and quarter-wave plates \hat{W}

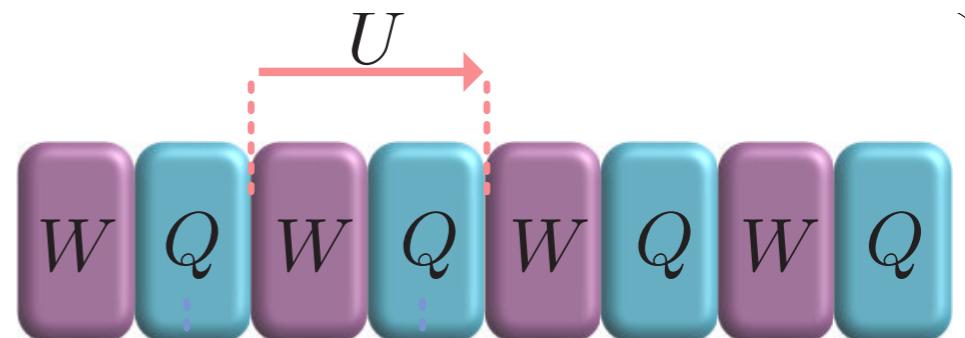
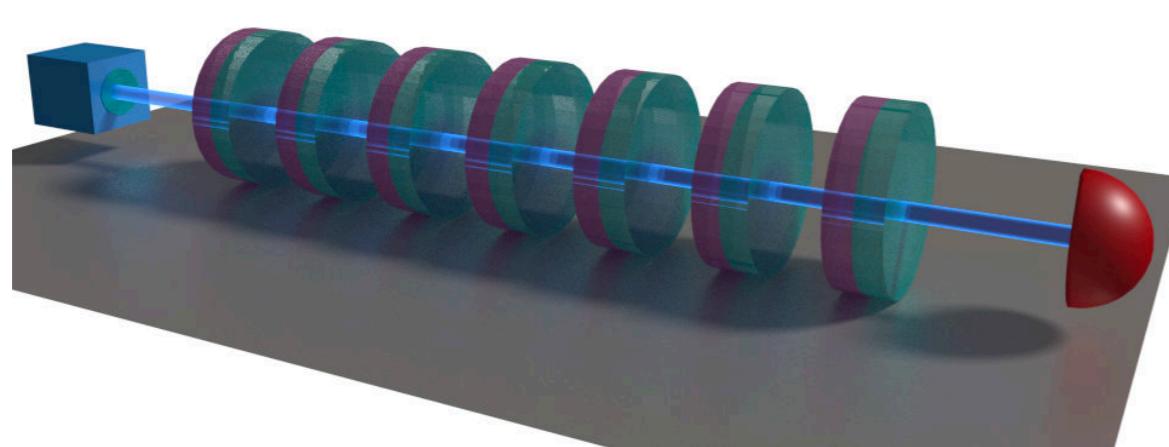
$$\hat{W} = \frac{1}{2} \begin{pmatrix} 1 & -i \\ -i & 1 \end{pmatrix}$$

discrete-time QW	Twisted photons
walker's position	OAM (m)
coin state (\uparrow/\downarrow)	polarization ($\circlearrowleft/\circlearrowright$)
spin rotation	W-plate
conditional displacement	Q-plate
time	\hat{z}



[Cardano et al., Science Advances (2015)]

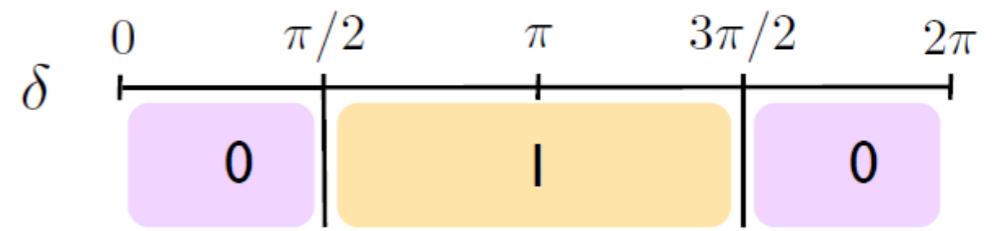
Discrete-Time Quantum Walk



- Periodic evolution: may be treated via Floquet theory
- One-step evolution operator $U \rightarrow H_{\text{eff}} \equiv \frac{i}{T} \log U$
- In momentum space: $H_{\text{eff}}(k) = E_k \hat{\mathbf{n}}_k \cdot \boldsymbol{\sigma}$
- The spectrum of H_{eff} is 2π -periodic (quasi-energies E_k)
- T+C+S symmetries: BDI class \rightarrow same invariant as the static SSH model

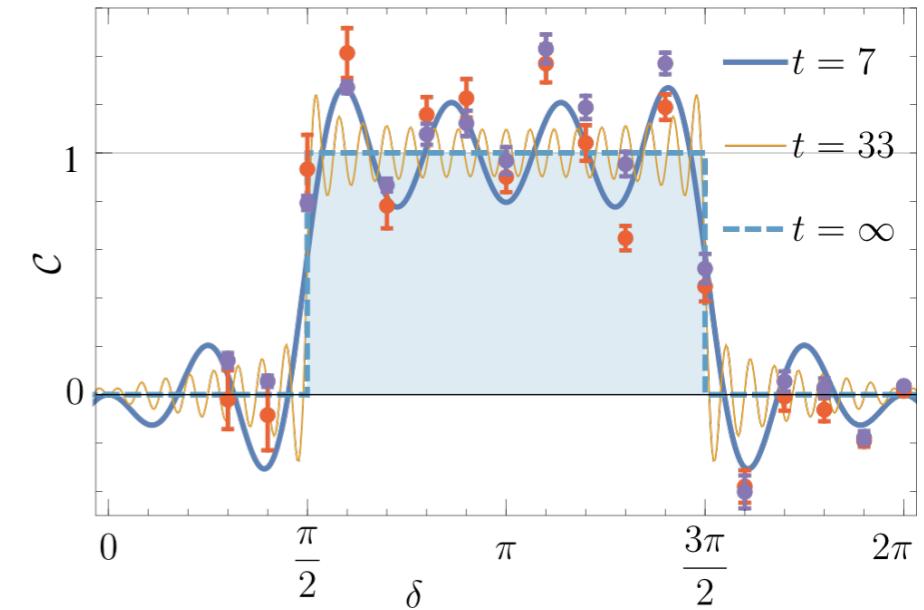
Detecting the invariant

- Winding: $\mathcal{W} = \oint \frac{dk}{2\pi} (\mathbf{n} \times \partial_k \mathbf{n})_z$



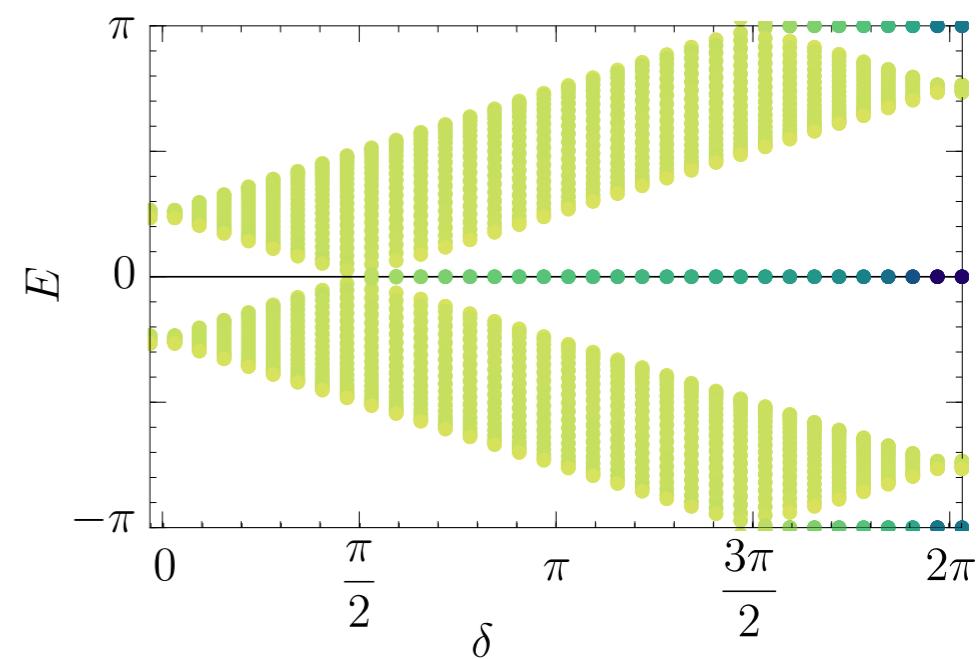
- Experimental measurement of the MCD after 7 timesteps of the DTQW with twisted photons:

(●/○): different initial polarizations



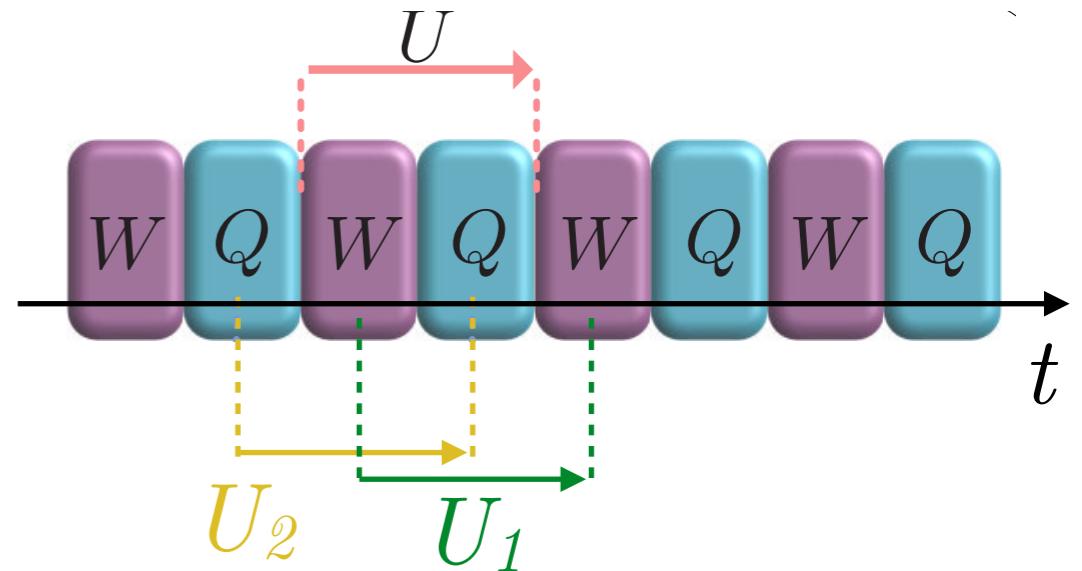
- Check bulk-boundary correspondence
- Spectrum with edges:

darker colors:
“edgier” states



Timeframes

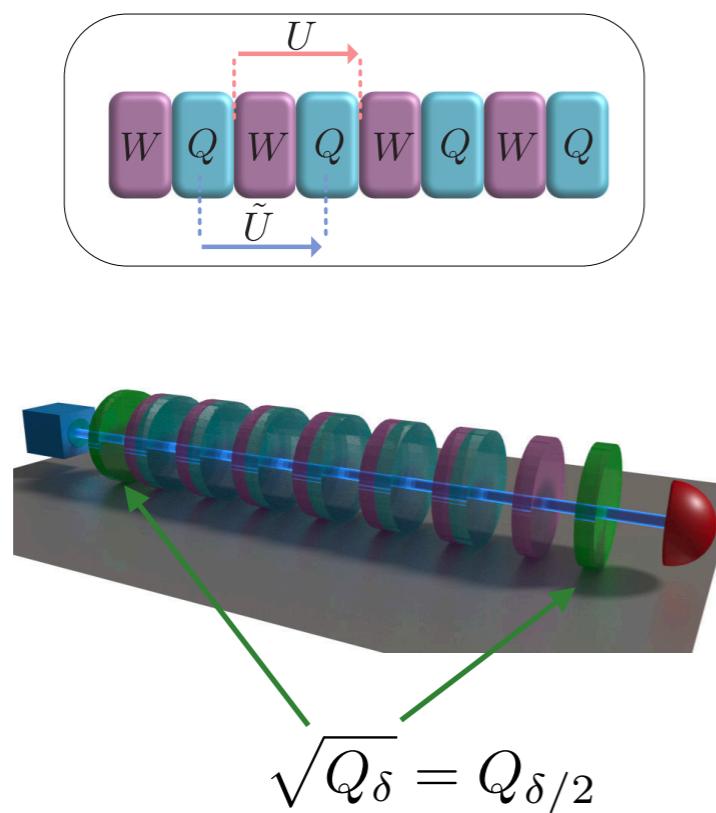
- Different initial t_0 lead to different U
- Eigenvalues of H_{eff} don't depend on t_0
- Eigenstates instead do! And so does the winding: $\mathcal{W} = \mathcal{W}_1 \neq \mathcal{W}_2$
- Timeframes invariant under time-reflection (U_1 and U_2) are special
- # of 0-energy edge states: $C_0 = (\mathcal{W}_1 + \mathcal{W}_2)/2$
- # of π -energy edge states: $C_\pi = (\mathcal{W}_1 - \mathcal{W}_2)/2$



[Asboth and Obuse, PRB (2013)]

Recovering the bulk-boundary correspondence

Measurement of the MCD in an alternative timeframe:



Spectrum (theory):

Measurements of C_0 and C_π :

Complete topological characterization
of a Floquet topological insulator

